

## Continuity equation

We consider the liquid, which moves in a space.

We use cartesian coordinates and consider, that any point of space can be described by three coordinates  $(x, y, z)$ .

The density of the liquid is a scalar function from coordinates. Let's denote density at the point  $(x, y, z)$  at the time moment  $t$  as  $\rho(x, y, z, t)$ .

Velocity of the liquid in the point of space is the vector-function from coordinates and time. It can be

described by its 3 values,  $u(x, y, z, t)$ ,  $v(x, y, z, t)$ ,  $w(x, y, z, t)$ , which are its projections to the axes at the moment  $t$ .

We assume, that functions  $\rho(x, y, z, t)$ ,  $u(x, y, z, t)$ ,  $v(x, y, z, t)$ ,  $w(x, y, z, t)$  are defined in a whole space and are differentiable there, and their derivative is continuous.

Let's consider some parallelepiped in the volume, with sides parallel to the axes.

Let size of sides are  $\Delta x, \Delta y, \Delta z$ , and coordinates of center of parallelepiped are  $(x_0, y_0, z_0)$ .

We divide the bottom of parallelepiped into  $N \times N$  equal squares with sides  $\delta x$  and  $\delta y$ .

Let's observe the system at some moment  $t_0$  for a small period of time  $\Delta t$ .

If sides of squares are small enough we may consider, that density and velocity doesn't change sufficiently in the square  $(x, y, x + \delta x, y + \delta y)$  in the small period  $\Delta t$ .

Mass of liquid, which comes through the square  $(x, y, x + \delta x, y + \delta y)$  is

$$u_z(x, y, z_0) \rho(x, y, z_0) \Delta t \delta x \delta y$$

Total mass, which comes through the bottom of parallelepiped is sum of mass coming through all squares

$$\Delta m_z = \sum_{i=1}^N \sum_{j=1}^N u_z(x_i, y_j, z_0, t_0) \Delta t \delta x \delta y$$

$$\Delta m_z = \Delta t \int_{x_0}^{x_0 + \Delta x} \int_{y_0}^{y_0 + \Delta y} \rho(x, y, z_0, t_0) u(x, y, z_0, t_0) dx dy$$

When  $\delta x, \delta y \rightarrow 0$ , sum mutates to the integral:

Now, we may introduce new function:

$$q_z\left(x + \frac{\Delta x}{2}, y + \frac{\Delta y}{2}, z, t\right) = \frac{1}{\Delta x \Delta y} \frac{\Delta m_z}{\Delta t} = \frac{1}{\Delta x \Delta y} \int_x^{x + \Delta x} \int_y^{y + \Delta y} \rho(x, y, z, t) u(x, y, z, t) dx dy$$

In the middle of the square it has a value, which shows, what mass of liquid goes through the square in  $z$  direction per unit of time per unit of square.

One may call it "flow in direction  $z$ ", or "speed of mass changing in direction  $z$ " or whatever his imagination allows him. We will call it " $q_z$ -function".

$$\text{In the similar way can be defined functions } q_x\left(x, y + \frac{\Delta y}{2}, z + \frac{\Delta z}{2}, t\right) \text{ and } q_y\left(x + \frac{\Delta x}{2}, y, z + \frac{\Delta z}{2}, t\right):$$

$$q_x\left(x, y + \frac{\Delta y}{2}, z + \frac{\Delta z}{2}, t\right) = \frac{1}{\Delta y \Delta z} \frac{\Delta m_x}{\Delta t} = \frac{1}{\Delta y \Delta z} \int_y^{y + \Delta y} \int_z^{z + \Delta z} \rho(x, y, z, t) u(x, y, z, t) dy dz$$

$$q_y\left(x + \frac{\Delta x}{2}, y, z + \frac{\Delta z}{2}, t\right) = \frac{1}{\Delta x \Delta z} \frac{\Delta m_y}{\Delta t} = \frac{1}{\Delta x \Delta z} \int_x^{x + \Delta x} \int_z^{z + \Delta z} \rho(x, y, z, t) v(x, y, z, t) dx dz$$

To make formulae shorter, let's denote

$$\bar{x} = x_0 + \frac{\Delta x}{2}, \quad \bar{y} = y_0 + \frac{\Delta y}{2}, \quad \bar{z} = z_0 + \frac{\Delta z}{2}$$

Then the mass of liquid, which comes through the bottom side of parallelipid during the time  $\Delta t$  is  $q_z(\bar{x}, \bar{y}, z_0, t_0) \Delta x \Delta y \Delta t$

The mass of liquid, which leaves through the top side of parallelipid during the time  $\Delta t$  is  $q_z(\bar{x}, \bar{y}, z_0 + \Delta z, t_0) \Delta x \Delta y \Delta t$

So, in parallelipid lefts  $\Delta x \Delta y \Delta t (q_z(\bar{x}, \bar{y}, z_0, t_0) - q_z(\bar{x}, \bar{y}, z_0 + \Delta z, t_0))$  mass of liquid, which comes through top or bottom side.

For the right-left and front-backward sides all is analogical.  
So, total change of mass in the parallelipid is:

$$\begin{aligned} \Delta m = & \Delta y \Delta z \Delta t (q_x(x_0, \bar{y}, \bar{z}, t_0) - q_x(x_0 + \Delta x, \bar{y}, \bar{z}, t_0)) + \\ & \Delta x \Delta z \Delta t (q_y(\bar{x}, y_0, \bar{z}, t_0) - q_y(\bar{x}, y_0 + \Delta y, \bar{z}, t_0)) + \\ & \Delta x \Delta y \Delta t (q_z(\bar{x}, \bar{y}, z_0, t_0) - q_z(\bar{x}, \bar{y}, z_0 + \Delta z, t_0)) \end{aligned}$$

Let's introduce jet another function: the mean density over the parallelipid

$$\bar{\rho}\left(x + \frac{\Delta x}{2}, y + \frac{\Delta y}{2}, z + \frac{\Delta z}{2}, t\right) = \bar{\rho}(\bar{x}, \bar{y}, \bar{z}, t) = \int_x^{x+\Delta x} \int_y^{y+\Delta y} \int_z^{z+\Delta z} \rho(x, y, z, t) dx dy dz$$

Then, change in mass is connected to change of mean density:

$$\begin{aligned} \Delta m &= m(t_0 + \Delta t) - m(t_0) = \bar{\rho}(\bar{x}, \bar{y}, \bar{z}, t_0 + \Delta t) V - \bar{\rho}(\bar{x}, \bar{y}, \bar{z}, t_0) V \\ \Delta m &= (\bar{\rho}(\bar{x}, \bar{y}, \bar{z}, t_0 + \Delta t) - \bar{\rho}(\bar{x}, \bar{y}, \bar{z}, t_0)) \Delta x \Delta y \Delta z \end{aligned}$$

And we have an equation:

$$\begin{aligned} & \Delta y \Delta z \Delta t (q_x(x_0, \bar{y}, \bar{z}, t_0) - q_x(x_0 + \Delta x, \bar{y}, \bar{z}, t_0)) + \\ & \Delta x \Delta z \Delta t (q_y(\bar{x}, y_0, \bar{z}, t_0) - q_y(\bar{x}, y_0 + \Delta y, \bar{z}, t_0)) + \\ & \Delta x \Delta y \Delta t (q_z(\bar{x}, \bar{y}, z_0, t_0) - q_z(\bar{x}, \bar{y}, z_0 + \Delta z, t_0)) = \\ & \quad = (\bar{\rho}(\bar{x}, \bar{y}, \bar{z}, t_0 + \Delta t) - \bar{\rho}(\bar{x}, \bar{y}, \bar{z}, t_0)) \Delta x \Delta y \Delta z \end{aligned}$$

If we divide both parts by  $\Delta x \Delta y \Delta z \Delta t$  we obtain

$$\frac{q_x(x_0, \bar{y}, \bar{z}, t_0) - q_x(x_0 + \Delta x, \bar{y}, \bar{z}, t_0)}{\Delta x} + \frac{q_y(\bar{x}, y_0, \bar{z}, t_0) - q_y(\bar{x}, y_0 + \Delta y, \bar{z}, t_0)}{\Delta y} + \frac{q_z(\bar{x}, \bar{y}, z_0, t_0) - q_z(\bar{x}, \bar{y}, z_0 + \Delta z, t_0)}{\Delta z} = \frac{\bar{\rho}(\bar{x}, \bar{y}, \bar{z}, t_0 + \Delta t) - \bar{\rho}(\bar{x}, \bar{y}, \bar{z}, t_0)}{\Delta t}$$

And now, if we take limit from that equation when  $\Delta x, \Delta y, \Delta z, \Delta t \rightarrow 0$   
We obtain partial differential equation:

$$-\frac{\partial q_x}{\partial x} - \frac{\partial q_y}{\partial y} - \frac{\partial q_z}{\partial z} = \frac{\partial \rho}{\partial t}$$

The only thing left is to take a look at the definition of functions  $q$ , when  $\Delta x, \Delta y, \Delta z, \Delta t \rightarrow 0$

$$q_y(x + \frac{\Delta x}{2}, y, z + \frac{\Delta z}{2}, t) = \frac{1}{\Delta x \Delta z} \int_x^{x+\Delta x} \int_z^{z+\Delta z} \rho(x, y, z, t) u(x, y, z, t) dx dz$$

when  $\Delta x, \Delta y, \Delta z, \Delta t$  are small, values of density and velocity do not change sufficiently.

So, in the limit case we may consider  $\rho(x, y, z, t) = \rho(x_0, y_0, z_0, t_0)$ ,

$$q_y(x + \frac{\Delta x}{2}, y, z + \frac{\Delta z}{2}, t) = q_y(x_0, y_0, z_0, t_0)$$

$$u(x, y, z, t) = u(x_0, y_0, z_0, t_0)$$

We have

$$q_y(x_0, y_0, z_0, t_0) = \frac{1}{\Delta x \Delta z} \int_x^{x+\Delta x} \int_z^{z+\Delta z} \rho(x_0, y_0, z_0, t_0) u(x_0, y_0, z_0, t_0) dx dz = \rho(x_0, y_0, z_0, t_0) u_y(x_0, y_0, z_0, t_0) \frac{\int_x^{x+\Delta x} \int_z^{z+\Delta z} dx dz}{\Delta x \Delta z} = \rho(x_0, y_0, z_0, t_0) u_y(x_0, y_0, z_0, t_0)$$

The same is for other components:

$$q_x(x_0, y_0, z_0, t_0) = \rho(x_0, y_0, z_0, t_0) u_x(x_0, y_0, z_0, t_0)$$

$$q_z(x_0, y_0, z_0, t_0) = \rho(x_0, y_0, z_0, t_0) u_z(x_0, y_0, z_0, t_0)$$

So, finally partial differential equation can be rewritten as

$$-\frac{\partial(\rho u_x)}{\partial x} - \frac{\partial(\rho u_y)}{\partial y} - \frac{\partial(\rho u_z)}{\partial z} = \frac{\partial \rho}{\partial t}$$

$$\text{or } \frac{\partial \rho}{\partial t} + \frac{\partial(\rho u_x)}{\partial x} + \frac{\partial(\rho u_y)}{\partial y} + \frac{\partial(\rho u_z)}{\partial z} = 0$$

$$\mathbf{u} = \begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix}$$

If we denote then equation may be rewritten in a more compact way:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{u} = 0$$

If the liquid is incompressible, then it has constant density. In that case equation becomes:

$$\nabla \cdot \mathbf{u} = 0$$